Some thoughts about modelling aggregation

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General properties of considered systems

- Pavel Krapivsky, Sidney Redner, Eli Ben-Naim. A kinetic view of statistical physics. Cambridge University Press, 2010.
- 📔 Фукс, Н. А. "Механика аэрозолей." 1955.
- Leyvraz Francois. Scaling theory and exactly solved models in the kinetics of irreversible aggregation. Physics Reports, 2003.

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Галкин В. А. "Уравнение Смолуховского", ФизМатЛит, 2001

General properties of considered systems

We consider system of chaotically moving particles which

- Collide inelastically
- Fill space homogeneously and uniformly
- Collide by pairs, triplets etc
- Differ in sizes

Aggregation and fragmentation processes can be described with Smoluchowski-type equations



- Aerosol dynamics
- Reversible polymerization
- Aggregation and fragmentation in planetary rings

General properties of considered systems



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General structure of models

In compact notations for the discrete case:

$$\begin{cases} \frac{d\mathbf{n}}{dt} = \mathbf{S}(\mathbf{n(t)}), \\ \mathbf{n(0)} = \mathbf{n_0} \end{cases}$$

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$$\mathbf{n}(t) = \begin{bmatrix} n_1(t) & n_2(t) & \dots & n_M(t) & \dots \end{bmatrix}$$
$$\mathbf{n}_0 = \begin{bmatrix} n_{1_0} & n_{2_0} & \dots & n_{M_0} & \dots \end{bmatrix},$$
and continuous

$$\begin{cases} \frac{\partial n(v,t)}{\partial t} = \mathbf{S}(n(t)), \\ n(v,t=0) = n_0(v) \end{cases}$$

Multi-component case:

$$\overline{v} = (v_1, v_2, \dots, v_d)$$

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Pairwise models

Continuous Smoluchowski equation

$$\frac{\partial n(v,t)}{dt} = \frac{1}{2} \int_0^v C(u,v-u)n(u,t)n(v-u,t)du - n(v,t) \int_0^\infty C(v,u)n(u,t)du.$$

Discrete Smoluchowski equation

$$\frac{dn_k(t)}{dt} = \frac{1}{2} \sum_{i+j=k} C_{i,j} n_i n_j - n_k \sum_{j\geq 1} C_{j,k} n_j , \ k = \overline{1,\infty} .$$

[Smoluchowski, M. von. Drei vortrage uber diffusion, brownsche bewegung und koagulation von kolloidteilchen. Zeitschrift fur Physik 17 (1916): 557-585]

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Pairwise models

• Irreversible coagulation with source and sink

$$\frac{dn_k(t)}{dt} = \frac{1}{2} \sum_{i+j=k} C_{i,j} n_i n_j - n_k \sum_{j=1}^N C_{j,k} n_j + J \delta_{k,1} , \ k = \overline{1,N} .$$

[Ball, Connaughton, Jones, Rajesh, Zaboronski, *Collective oscillations in irreversible coagulation driven by monomer inputs and large-cluster outputs*, PRL (2012)]

Aggregation-fragmentation equations in planetary rings

$$\frac{dn_k(t)}{dt} = \frac{1}{2} \sum_{i+j=k} C_{i,j} n_i n_j - (1+\lambda) n_k \sum_{j\geq 1} C_{j,k} n_j , \ k = \overline{2,\infty}.$$

$$\frac{dn_1}{dt} = -n_1 \sum_{j\geq 1} C_{1,j}n_j + \frac{\lambda}{2} \sum_{i,j\geq 2} C_{i,j}(i+j)n_in_j + \lambda n_1 \sum_{j\geq 2} jC_{1,j}n_j.$$

[Brilliantov, Krapivsky, Bodrova, ... Size distribution of particles in Saturn's rings from aggregation and fragmentation, PNAS (2015)]

Pairwise models

• Multicomponent Smoluchowski equation

$$\frac{\partial n(\overline{v},t)}{\partial t} = \frac{1}{2} \int_0^{v_1} \dots \int_0^{v_d} C(\overline{v} - \overline{u};\overline{u}) n(\overline{v} - \overline{u},t) n(\overline{u},t) du_1 \dots du_d - n(\overline{v},t) \int_0^{\infty} \dots \int_0^{\infty} C(\overline{u};\overline{v}) n(\overline{u},t) du_1 \dots du_d.$$

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Multi-particle kinetics

$$\begin{aligned} \frac{dn_k}{dt} &= S_k^{(2)}(\mathbf{n}) + S_k^{(3)}(\mathbf{n}), \qquad k = \overline{1, \infty}, \\ S_k^{(2)}(\mathbf{n}) &= \frac{1}{2} \sum_{i_1+i_2=k} C_{i_1,i_2}^{(2)} n_{i_1} n_{i_2} - n_k \sum_{i_1=1}^{\infty} C_{i_1,k}^{(2)} n_{i_1}, \\ S_k^{(3)}(\mathbf{n}) &= \frac{1}{6} \sum_{i_1+i_2+i_3=k} C_{i_1,i_2,i_3}^{(3)} n_{i_1} n_{i_2} n_{i_3} - \frac{n_k}{2} \sum_{i_1,i_2=1}^{\infty} C_{i_1,i_2,k}^{(3)} n_{i_1} n_{i_2}. \end{aligned}$$

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Low-rank decompositions

Matrix skeleton decompositon



and rule for evaluation of the elements:

$$X(i_1,i_2) = \sum_{\alpha=1}^{R} U(i_1,\alpha)V(i_2,\alpha)$$

- We store just $2nR \ll n^2$. It is good when $R \ll n$
- There are robust methods

Low-rank decompositions

Canonical polyadic



$$X(i_1, i_2, \ldots, i_d) = \sum_{\alpha=1}^R U_1(i_1, \alpha) U_2(i_2, \alpha) \ldots U_d(i_d, \alpha)$$

- Compression $O(N^d) \rightarrow O(dNR)$
- Evaluation of rank is NP-full

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Low-rank decompositions

Tensor train

$$X_{i_1,i_2,\ldots,i_d} = \sum_{\alpha_i=1}^R G_1(\alpha_0,i_1,\alpha_1)G_2(\alpha_1,i_2,\alpha_2)\ldots G_d(\alpha_{d-1},i_d,\alpha_d)$$



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- Compression $O(N^d)
 ightarrow O(dNR^2)$
- Complexity O(dNR³)
- Robust methods
- Robust operations

Predictor-corrector time-integration

$$\frac{n^{t+\frac{1}{2}}-n^{t}}{0.5\tau}=S^{(2)}(n^{t})+S^{(3)}(n^{t})+\ldots+S^{(D)}(n^{t})$$
$$\frac{n^{t+1}-n^{t}}{\tau}=S^{(2)}(n^{t+\frac{1}{2}})+S^{(3)}(n^{t+\frac{1}{2}})+\ldots+S^{(D)}(n^{t}),$$

Straight-forward complexity of evaluation of $S^{(D)}$ is $O(N^D)$. Even for two-component case it is too much!

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- TT-based algorithm $O(NDR^2 \log N)$
- CP-based algorithm $O(NDR \log N)$

Main idea for pairwise case

Re-write ${f S}$ as:

$$\frac{1}{2} \sum_{i=2}^{N} \sum_{j=2}^{N} C_{i,j}(i+j)n_in_j = \frac{1}{2} \sum_{i=2}^{N} in_i \sum_{j=2}^{N} C_{i,j}n_j + \frac{1}{2} \sum_{i=2}^{N} n_i \sum_{j=2}^{N} C_{i,j}jn_j = \\ = \begin{bmatrix} n_2 & n_3 & \dots & n_N \end{bmatrix} \times \begin{bmatrix} C_{2,2} & C_{2,3} & \dots & C_{2,N} \\ C_{3,2} & \dots & \dots & C_{3,N} \\ \dots & \dots & \dots & \dots \\ C_{N,2} & C_{N,3} & \dots & C_{N,N} \end{bmatrix} \times \begin{bmatrix} 2n_2 \\ 3n_3 \\ \dots \\ Nn_N \end{bmatrix} \approx \\ \approx \begin{bmatrix} n_2 & n_3 & \dots & n_N \end{bmatrix} \times UV^T \times \begin{bmatrix} 2n_2 \\ 3n_3 \\ \dots \\ Nn_N \end{bmatrix}$$

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Main idea for pairwise case

$$\sum_{i=1}^{k-1} C_{i,j} n_i n_{k-i} \approx \sum_{i=1}^{k-1} \sum_{\alpha=1}^{R} U_{\alpha}(i) V_{\alpha}(k-i) n_i n_{k-i} =$$
$$= \sum_{i=1}^{k-1} \sum_{\alpha=1}^{R} \widehat{U_{\alpha}}(i) \widehat{V_{\alpha}}(k-i)$$
$$\widehat{U_{\alpha}}(i) \equiv U_{\alpha}(i) n_i; \ \widehat{V_{\alpha}}(i) \equiv V_{\alpha}(i) n_i.$$

Hence complexity of time-step becomes $O(NR \log N)$ operations. Then we need estimetes for R.

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Oscillations in aggregation shattering processes



 $a = 0.95, \lambda = 0.005$



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Oscillations in aggregation with source and sink

$$\frac{dn_k(t)}{dt} = \frac{1}{2} \sum_{i+j=k} C_{i,j} n_i n_j - n_k \sum_{j=1}^N C_{j,k} n_j + J \delta_{k,1} , \ k = \overline{1, N} .$$

There is trivial scaling rule allowing to normalize sink:

$$n_i(t)|_J = n_i\left(t\cdot\sqrt{J}\right)|_{J=1}\cdot\sqrt{J}.$$



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Advection-coagulation in 2D spatial case



$$\frac{\partial f(t, x, y, v)}{\partial t} + \nabla \cdot (c(x, y, v)f(t, x, y, v)) =$$

$$= \frac{1}{2} \int_{0}^{v} K(u, v - u)f(t, x, y, u)f(t, x, y, v - u)du -$$

$$-f(t, x, y, v) \int_{0}^{\infty} K(u, v)f(t, x, y, u)du,$$

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Issues in progress for advection-coaulation

- Add realistic evaluation of velocities c(x, y, v)
- Introduce some theoretical anlysis for convergence and stability
- Incorporate use of ready-made CFD packages and platforms

We are very interested in discussions and advice!

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Main idea for milti-component case

Represent both kernel and soluton in TT-format and use TT-operations

$$\frac{\partial n(\overline{v},t)}{\partial t} = \frac{1}{2} \int_0^{v_1} \dots \int_0^{v_d} C(\overline{v} - \overline{u};\overline{u}) n(\overline{v} - \overline{u},t) n(\overline{u},t) du_1 \dots du_d - n(\overline{v},t) \int_0^{\infty} \dots \int_0^{\infty} C(\overline{u};\overline{v}) n(\overline{u},t) du_1 \dots du_d.$$

• After some technical preparations we obtain complexity of time-step $O(dN \log NR^3 + d^2NR^6)$ instead of $O(N^{2d})$

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Multicomponent case performance vs Monte Carlo





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Numerical results for triple collisions



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Speedup in parallel¹

| Number of CPU-cores | time, sec | Speedup |
|---------------------|-----------|---------|
| 1 | 257.80 | 1.00 |
| 2 | 147.62 | 1.75 |
| 4 | 80.21 | 3.21 |
| 8 | 43.65 | 5.91 |
| 16 | 22.63 | 11.39 |
| 32 | 14.83 | 17.38 |
| 64 | 13.15 | 19.60 |
| 128 | 12.22 | 21.09 |

¹"Zhores" supercomputer of Skoltech, 100 Gb/s EDR, 24 CPU cores per node

Спасибо за внимание! Thank you for attention! 2

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²Our codes are open, ask if you interested